

## Volume of tubes in noncompact symmetric spaces

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**Abstract.** From the duality between symmetric spaces of compact type and symmetric spaces of noncompact type, we deduce an expression for the “infinitesimal change of volume function” of geodesic balls and tubes around certain submanifolds of a noncompact symmetric space of arbitrary rank. We also give a formula for the volume of a geodesic ball in a noncompact symmetric space.

### 0. Introduction

The theory of tubes and geodesic balls in symmetric spaces of rank one has been fully developed without making an explicit reference mention to the theory of roots of the symmetric space [AGV], [G].

One important fact that arises for compact symmetric spaces of rank one is that the isotropy representation acts transitively on the unit sphere of the tangent space of any point. Since that property is not satisfied for symmetric spaces of rank greater than one, we have to use the theory of restricted roots and orbital geometry of the linear action of the isotropy group of a compact symmetric space, to obtain an expression for the area of geodesic balls and tubes around a submanifold of that compact symmetric space [N-G].

The aim of this work is to use the duality between compact and noncompact symmetric spaces to show that some of the formulae obtained in

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*Mathematics Subject Classification:* 53C35, 53C21.

*Key words and phrases:* noncompact symmetric space, geodesic ball, restricted roots, tube, totally geodesic submanifold.

Work partially supported by the Spanish DGICYT grant No. PB94-0972 and the Mexican CONACyT grant No. 0329P-E.

[N-G] for geodesic balls and tubes in compact symmetric spaces of arbitrary rank are easily adaptable to the noncompact case.

Finally, if we restrict our results to symmetric spaces of rank one we can deduce the well-known formulae for geodesic balls and tubes in these spaces. For instance, the volume of a geodesic ball in an  $n$ -dimensional sphere of constant curvature  $\lambda = 1$  is given by

$$(1) \quad V_m^1(r) = \frac{2\pi^{\frac{n}{2}}}{\Gamma(\frac{n}{2})} \int_0^r \sin^{n-1}(t) \, dt,$$

and the volume of a geodesic ball in the hyperbolic space of constant curvature  $\lambda = -1$  is

$$(2) \quad V_m^2(r) = \frac{2\pi^{\frac{n}{2}}}{\Gamma(\frac{n}{2})} \int_0^r \sinh^{n-1}(t) \, dt.$$

So, the formulas for the case of sectional curvature  $\lambda < 0$  can be found by changing all of the trigonometric functions, which appear for the case of sectional curvature  $\lambda > 0$ , to hyperbolic functions. Similar remarks apply to all rank one symmetric spaces and, as we will see in this work, also for symmetric spaces of arbitrary rank.

To illustrate this program we will use the symmetric spaces of type **AIII** as examples. Therefore, the complex projective space (compact case) and the complex hyperbolic space (noncompact case) are included in these examples.

### 1. Duality for symmetric spaces

Let  $M_0 = G_0/K$  be a noncompact symmetric space and  $M = G/K$  a compact symmetric space dual to  $M_0$ . We denote by  $\mathfrak{g}_0$  the orthogonal involutive Lie algebra associated to  $M_0$ . Let

$$(3) \quad \mathfrak{g}_0 = \mathfrak{k} + \mathfrak{m}_0$$

be the Cartan decomposition of  $\mathfrak{g}_0$  ( $\mathfrak{k}$  is the Lie algebra of  $K$ ). If  $\mathfrak{g}$  denotes the complexification of  $\mathfrak{g}_0$ , then the compact real form  $\mathfrak{g}_c$  of  $\mathfrak{g}$ , given by

$$(4) \quad \mathfrak{g}_c = \mathfrak{k} + \mathfrak{m} \quad \text{where } \mathfrak{m} = i\mathfrak{m}_0, \quad (i = \sqrt{-1}),$$

is the orthogonal Lie algebra associated with  $M$ .

Let  $h_{m_0}$  denote a maximal abelian subspace of  $m_0$  and let  $h_0$  be any maximal abelian subalgebra of  $g_0$  containing  $h_{m_0}$ . We denote by  $h$  the subspace of  $g$  generated by  $h_0$  and  $h_m = ih_{m_0}$  denotes a maximal abelian subspace of  $m$ .

We give  $M_0$  the unique  $G_0$ -invariant Riemannian structure induced by the restriction of the Killing form  $B$  of  $g$  to  $m_0 \times m_0$ . Similarly, the space  $M$  will be given the unique  $G$ -invariant Riemannian structure induced by the restriction of  $-B$  of  $g$  to  $m \times m$ .

Let  $\Delta$  denote the set of all nonzero roots of  $g$  with respect to  $h$ , and let  $\Delta_{m_0}$  denote the set of those roots which do not vanish identically on  $h_{m_0}$ . Each  $\alpha \in \Delta$  is a  $\mathbb{C}$ -linear function on  $h$  which is real-valued on  $h_{m_0}$ ; therefore, we get in this way an ordering of  $\Delta_{m_0}$ . Let  $\Delta^+$  denote the subset of positive roots in  $\Delta_{m_0}$ .

The sectional curvature of the symmetric spaces  $M_0$  and  $M$  is given, by means of the roots in  $\Delta^+$ , as follows:

Let  $H_0$  and  $X_0$  be two unit vectors in  $h_{m_0}$  and  $m_0$ , respectively; the sectional curvature of  $M_0$  is

$$(5) \quad \begin{aligned} K_{H_0X_0} &= \langle R(H_0, X_0)H_0, X_0 \rangle \\ &= \langle -[H_0, [H_0, X_0]], X_0 \rangle = -\alpha^2(H_0), \quad \alpha \in \Delta^+. \end{aligned}$$

Therefore, since  $\alpha(H_0) \in \mathbb{R}$ ,  $K_{H_0X_0} < 0$ .

Now, let  $H = iH_0$  and  $X = iX_0$  be two unit vectors in  $h_m$  and  $m$ , respectively; the sectional curvature of  $M$  is

$$(6) \quad K_{HX} = \langle R(H, X)H, X \rangle = -\alpha^2(H) = -\alpha^2(iH_0) = \alpha^2(H_0) > 0.$$

On the other hand, if we consider the linear actions

$$\text{Ad} : K \times m_0 \longrightarrow m_0 \quad \text{and} \quad \text{Ad} : K \times m \longrightarrow m,$$

it can be shown that the two actions have the same principal orbit type. Moreover, if we denote by  $K/K_{H_0}$  this principal orbit type, ( $K_{H_0}$  is the isotropy subgroup of  $K$ ), the volume of the principal orbits  $K(H_0) = \bigcup_{k \in K} \text{Ad}(k)H_0$ , ( $H_0 \in m_0$ ), and  $K(H) = \bigcup_{k \in K} \text{Ad}(k)H$ , ( $H \in m$ ) is  $([H-H], [H-L])$ :

$$(7) \quad v^2(K(H_0)) = c^2 \prod_{j=q+1}^n \alpha_j^2(H_0) = c^2 \prod_{j=q+1}^n \alpha_j^2(H) = v^2(K(H));$$

where  $c$  denotes the constant given by the volume of  $K/K_{H_0}$ .

## 2. The volume function of a geodesic ball in a noncompact symmetric space

We begin this section by giving an expression for the infinitesimal change of the volume function from the roots in  $\Delta^+$ . The infinitesimal change of the volume function appears in the volume formula of a tube and it is defined as follows: given a submanifold  $P$  of  $M_0$ , the infinitesimal change of the volume function  $\vartheta_{H_0}(t)$  of  $P$  in the direction of a unit vector  $H_0$ , normal to  $P$ , is given by  $\omega(\xi(t)) = \vartheta_{H_0}(t) dH_0 \wedge dP \wedge dt$ , where  $\xi(t)$  is a geodesic of  $M$  satisfying  $\xi(0) = p \in P$  and  $\xi'(0) = H_0$ ,  $\omega$  is the Riemannian volume element of  $M_0$ ,  $dP$  that of  $P$ , and  $dH_0$  that of the unit sphere in the normal space  $T_p P^\perp$ . This definition makes sense when  $P$  is a point  $p \in M_0$  and  $H_0$  is a unit vector in  $m_0$ ; in this case a geodesic ball (tube around a point) is considered and we obtain the following expression for the infinitesimal change of the volume function:

**Proposition 2.1.** *Let  $M_0$  be a noncompact symmetric space and  $\xi(t)$  a geodesic in  $M_0$  such that  $\xi(0) = p$  and  $\xi'(0) = H_0$  ( $H_0$  a unit vector in a maximal abelian subspace  $h_{m_0}$  of  $m_0$ ). Then, the infinitesimal change of the volume function  $\vartheta_{H_0}(t)$  is given by*

$$(8) \quad \vartheta_{H_0}(t) = t^{-(n-q)} \prod_{j=q+1}^n \frac{\sinh(t\alpha_j(H_0))}{\alpha_j(H_0)}.$$

where  $q$  is the rank of  $M_0$ , (i.e.  $\dim(h_{m_0})$ ), and  $\alpha_j$  varies in  $\Delta^+$ , where without loss of generality we assume that  $0 = \alpha_1(H_0) = \dots = \alpha_q(H_0) < \alpha_{q+1}(H_0) \leq \dots \leq \alpha_n(H_0)$ .

PROOF. By analogy with the compact case (see [N-G]), we omit the explicit proof of the proposition. The main difference between the two cases is that for noncompact symmetric spaces the sectional curvature is negative and the Riccati differential equations for the principal curvature functions are:

$$(9) \quad \begin{aligned} k'_i(t) &= k_i(t)^2, & i &= 2, \dots, q; \\ k'_j(t) &= k_j(t)^2 - \alpha_j(H_0)^2, & j &= q+1, \dots, n. \end{aligned}$$

Now, starting with the expression (8) for the infinitesimal change of volume function, our purpose is to simplify the integral over the unit sphere

which appears in the next identity for the volume function of a geodesic ball:

$$(10) \quad V_p^{M_0}(r) = \int_0^r A_p^{M_0}(t) \, dt = \int_0^r \int_{S^{n-1}(1)} t^{n-1} \vartheta_{H_0}(t) \, dH_0 \, dt.$$

In order to achieve this, we consider the next two properties, which are satisfied for compact and noncompact symmetric spaces.

**Proposition 2.2.** *The infinitesimal change of volume function  $\vartheta_{H_0}(t)$  does not depend on the vector chosen in the orbit  $\bigcup_{k \in K} Ad(k)H_0$ .*

**Proposition 2.3.** *Let  $S^{n-1}(1)$  and  $S^{q-1}(1)$  be the unit spheres in  $m_0$  and  $h_{m_0}$ , respectively. Given a Weyl chamber  $D_0$  in  $h_{m_0}$  we consider  $C_0 = \bar{D}_0 \cap S^{q-1}(1)$ , where  $\bar{D}_0$  is the closure of  $D_0$ ; then,*

$$(11) \quad S^{n-1}(1) = \bigcup_{k \in K} Ad(k)C_0.$$

(See [N-G] for the proofs of both propositions.)

**Theorem 2.4.** *The volume of a geodesic ball in the noncompact symmetric space  $M_0$  is given by*

$$(12) \quad \begin{aligned} V_p^{M_0}(r) &= \int_0^r A_p^{M_0}(t) \, dt \quad \text{where} \\ A_p^{M_0}(t) &= ct^{q-1} \int_{C_0} \prod_{j=q+1}^n \sinh(t\alpha_j(H_0)) \, dH_0 \end{aligned}$$

where  $c$  denotes the constant given by the volume of the principal orbit  $K/K_{H_0}$ .

PROOF. From Propositions 2.2 and 2.3, we can express the area function of the geodesic ball as

$$(13) \quad A_p^{M_0}(t) = t^{n-1} \int_{C_0} v(H_0) \vartheta_{H_0}(t) \, dH_0,$$

where  $v(H_0)$  is the volume of the orbit  $\bigcup_{k \in K} Ad(k)H_0$ .

Now, replacing  $\vartheta_{H_0}(t)$  and  $v(H_0)$  in (13), and having in mind (8) and (7), respectively, we obtain the desired result.  $\square$

**3. Example: Symmetric spaces of type AIII**

Now, we shall consider the noncompact symmetric space  $M_0 = SU(p, q)/S(U_p \times U_q)$ , and we specialize the expression (12) for this space. Here,

$$(14) \quad m = \left\{ \left( \begin{array}{cc} 0 & iZ \\ -i\bar{X}t & 0 \end{array} \right) : Z \text{ complex } p \times q \text{ matrix} \right\};$$

and

$$(15) \quad h_{m_0} = \sum_{j=1}^q i\mathbb{R}(E_{j,p+j} - E_{p+j,j})$$

where  $E_{jk}$  denotes the matrix  $(\delta_{aj}\delta_{bk})$ .

Let  $H_0$  be a unit vector in  $h_{m_0}$  which we identify with  $(x_1, \dots, x_q)$ , such that  $x_1^2 + \dots + x_q^2 = 1$ . Then, from (5), the eigenvalues of  $R_{H_0}$  are: 0 with multiplicity  $q$ ,  $-x_j^2$  with multiplicity  $2(p - q)$ ,  $-4x_j^2$  with multiplicity 1 and  $-(x_j - x_k)^2$ ,  $-(x_j + x_k)^2$  with multiplicity 2; for  $1 \leq j < k \leq q$ .

A Weyl chamber in  $h_{m_0}$  is given by  $D_0 = \{(x_1, \dots, x_q) : x_1 > \dots > x_q > 0\}$ .

From (12) we have that

$$(16) \quad \begin{aligned} A_p^{M_0}(t) &= ct^{q-1} \int_{C_0} \prod_{j=1}^{2(p-q)} \sinh(tx_j) \prod_{j=1}^q \sinh(2tx_j) \\ &\quad \times \prod_{\substack{j=1 \\ j < k}}^q \sinh^2(t(x_k - x_j)) \sinh^2(t(x_k + x_j)) \, dx_1 \dots dx_q. \end{aligned}$$

*Remark 1.* For  $q = 1$ ,  $M_0$  is the complex hyperbolic space and the area function of the geodesic ball is, [G],

$$(17) \quad A_p^{M_0}(t) = c \sinh^{2(p-1)}(t) \sinh(2t) = 2c \sinh^{2p-1}(t) \cosh(t).$$

*Remark 2.* The area function  $A_p^M(t)$  for the compact symmetric space dual to  $M_0$  is given by changing in (16) all the hyperbolic functions to trigonometric functions.

To end this section we plot the volume of a geodesic ball, up to the same constant factor, in the compact and noncompact symmetric spaces

of type **AIII** from 0 to the injectivity radius of the compact symmetric space, which is  $\pi/2$ , [N-G].

$$SU(4)/S(U(2) \times U(2))$$

$$SU(2, 2)/S(U(2) \times U(2))$$

#### 4. Tubes in noncompact symmetric spaces

In the preceding section we have computed the volume of a geodesic ball in a noncompact symmetric space  $M_0$  as an adaptation of the volume function of geodesic balls in the compact symmetric space  $M$  dual to  $M_0$ . In this section we consider tubes around totally geodesic submanifolds compatible with  $M_0$ .

Let  $P_0$  be a totally geodesic submanifold in  $M_0$ ; then ([He] and [W]),  $P_0 = U_0/L$  is a noncompact symmetric space. Let  $\mu_0$  denote the Lie algebra of  $U_0$  and  $\mu_0 = \rho_0 + l$  its canonical decomposition.

In order to get the compatibility condition of  $P_0$  with  $M_0$  we assume that for all vectors  $H_0$  in  $\rho_0^\perp$  the curvature operator satisfies the condition of preserving the curvature  $R_{H_0}(\rho_0) \subset \rho_0$  and  $R_{H_0}(\rho_0^\perp) \subset \rho_0^\perp$ , (the first inclusion is immediate because  $\rho_0$  is a Lie triple system).

Let  $P$  denote the compact symmetric space dual to  $P_0$ , which is, in its turn, a totally geodesic submanifold compatible with  $M$ .

In [N-G] we have computed the infinitesimal change of volume function of  $P$ ; therefore, as we have seen in Proposition 2.1 for geodesic balls, the infinitesimal change of the volume function of  $P_0$  is given by changing all of the trigonometric functions, which appear in the infinitesimal change of the volume function of  $P$ , to hyperbolic functions.

However, although we have an expression for the infinitesimal change of the volume function of  $P_0$ , it is not possible to obtain a finite value for volume of a tube around  $P_0$  because  $P_0$  is a noncompact submanifold of  $M_0$ .

*Acknowledgement.* We wish to express our gratitude to Dr. A. SANCHEZ-VALENZUELA for his fruitful comments about symmetric spaces and for making possible the visit of the first author to the CIMAT (Guanajuato, Mexico).

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(Received July 10, 1997; accepted March 17, 1998)